

On the Use of the DuFort-Frankel Finite-Difference Approximation for Simulation of Diffusion in Geophysical Fluids¹

JAMES E. OVERLAND

Dept. of Meteorology and Oceanography, New York University, Bronx, N. Y. 10453

16 January 1973 and 16 April 1973

The DuFort-Frankel scheme for numerical simulation of diffusion does not necessarily conserve scalar properties. Consider the temperature equation commonly

used in the Bénard cell problem:

$$\frac{T_{ij}^{N+1} - T_{ij}^{N-1}}{2\Delta t} = -J^*(\psi^N, T^N) + \frac{K}{(\Delta S)^2} \delta^N, \quad (1)$$

¹ Contribution No. 135, Geophysical Sciences Laboratory, New York University.

where J^* is a conservative Jacobian operator, ΔS the

space discretization, Δt the time step, K the thermal diffusivity, and δ a diffusion operator. Two numerically stable diffusion operators in common use are forward differencing with a one time step lag, i.e.,

$$\delta_F^N = T_{i+1,j}^{N-1} + T_{i-1,j}^{N-1} + T_{i,j+1}^{N-1} + T_{i,j-1}^{N-1} - 4T_{ij}^{N-1}, \quad (2)$$

and the DuFort-Frankel operator

$$\delta_{DF}^N = T_{i+1,j}^N + T_{i-1,j}^N + T_{i,j+1}^N + T_{i,j-1}^N - 2T_{ij}^{N+1} - 2T_{ij}^{N-1}. \quad (3)$$

Subscripts i and j refer to space points, while the superscript refers to time level.

Further, consider adjacent boxes A and B, with all sides insulated except for the interface. For the forward-difference scheme, the diffusive flux from A to B per time step is given by

$$\gamma [T_A^{N-1} - T_B^{N-1}], \quad (4)$$

with

$$\gamma \equiv \frac{2K\Delta t}{(\Delta S)^2}. \quad (5)$$

By symmetry the flux from B to A per time step is equal and opposite, which conserves the scalar T . For the DuFort-Frankel scheme, the flux from A to B is given by

$$\gamma \left[T_A^N - \left(\frac{T_B^{N+1} + T_B^{N-1}}{2} \right) \right]; \quad (6)$$

this is not necessarily the same magnitude as the flux

from B to A, which is given by

$$\gamma \left[T_B^N - \left(\frac{T_A^{N+1} + T_A^{N-1}}{2} \right) \right]. \quad (7)$$

In certain geophysical problems the non-conservative nature of the DuFort-Frankel scheme may be restrictive. Salinity intrusion in an estuary has several features in common with the Bénard problem, with the major differences being that a density source is maintained at a side boundary and that the height-to-length ratio is very small. Numerical experiments using the DuFort-Frankel scheme on the estuary problem consistently lost salt. The loss of salt was estimated to be 0.60% of the salt in the system per time iteration. It is felt that the DuFort-Frankel scheme did not perform well for this problem because the interior was far removed from the source-sink boundaries.

The use of the forward-difference scheme appears to be necessary when strict conservation of a scalar is required, although the required small time step may be prohibitive. The other recourse is the use of implicit methods, which can be handled easily in one dimension (Richtmyer and Morton, p. 199), but are much more cumbersome in large multi-dimensional systems. It is apparent that the inclusion of diffusion in numerical simulations is a non-trivial exercise.

Acknowledgments. This work was supported by the ARCA Foundation and by the National Sea Grant Program (Project Grant No. 1-36105). The author gratefully acknowledges the guidance of Dr. W. J. Pierson and Dr. J. Spar.

REFERENCE

Richtmyer, R. D., and K. W. Morton, 1967: *Difference Methods for Initial Value Problems*. New York, Interscience, 406 pp.