Recent Climate Variability in Antarctica from Satellite-Derived Temperature Data

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ABSTRACT

Recent Antarctic climate variability on month-to-month to interannual time scales is assessed through joint analysis of surface temperatures from satellite thermal infrared observations (T_{IR}) and passive microwave brightness temperatures (T_B). Although T_{IR} data are limited to clear-sky conditions and T_B data are a product of the temperature and emissivity of the upper ~1 m of snow, the two datasets share significant covariance. This covariance is largely explained by three empirical modes, which illustrate the spatial and temporal variability of Antarctic surface temperatures. The T_B variations are damped compared to T_{IR} variations, as determined by the period of the temperature forcing and the microwave emission depth; however, microwave emissivity does not vary significantly in time. Comparison of the temperature modes with Southern Hemisphere (SH) 500-hPa geopotential height anomalies demonstrates that Antarctic temperature anomalies are predominantly controlled by the principal patterns of SH atmospheric circulation. The leading surface temperature mode strongly correlates with the Southern Annular Mode (SAM) in geopotential height. The second temperature mode reflects the combined influences of the zonal wavenumber-3 and Pacific–South American (PSA) patterns in 500-hPa height on month-to-month time scales. ENSO variability projects onto this mode on interannual time scales, but is not by itself a good predictor of Antarctic temperature anomalies. The third temperature mode explains winter warming trends, which may be caused by blocking events, over a large region of the East Antarctic plateau. These results help to place recent climate changes in the context of Antarctica's background climate variability and will aid in the interpretation of ice core paleoclimate records.

1. Introduction

A number of studies have shown considerable interest in identifying and explaining Antarctic temperature trends over recent decades (Doran et al. 2002; Marshall 2002a; Thompson and Solomon 2002; Vaughan et al. 2001; van den Broeke 2000a). However, because the interannual variability of Antarctic climate is large, it is difficult to establish the significance of surface temperature trends from sparsely distributed weather stations on the continent (King 1994). Furthermore, relatively little is known about the spatial structure of surface temperature variations across Antarctica. Such knowledge would, for example, improve the interpretation of ice core paleoclimate records, which are usually obtained from locations that are remote from weather stations.

Two important influences on Antarctica's climate variability, the Southern Annular Mode (SAM) and the El Niño–Southern Oscillation (ENSO), have been discussed by several studies, and increased tendency for these circulation patterns to stay in a particular phase may be driving surface temperature trends in the Antarctic (Gillett and Thompson 2003; Bromwich et al. 2003; Ribera and Mann 2003; Thompson and Solomon 2002; Kwok and Comiso 2002). We are therefore motivated to pay particular attention to the influence of these atmospheric patterns on Antarctic surface temperature anomalies, which, in this study, are derived from passive microwave brightness temperature (T_B) and thermal infrared satellite observations (T_{IR}).

In previous work with these data, Schneider and Steig (2002, hereafter SS) presented a principal component analysis of T_B data and showed evidence for the SAM and ENSO-related signals in Antarctica. However, the T_B data, taken alone, can be complicated to interpret because of the effects of nonstationary microwave emissivity variations due to variations in snow characteris-

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tics, and occasional surface melt events. Kwok and Comiso (2002) examined newly available T_{IR} data, and also linked their variability to the SAM and ENSO. That study assumed, a priori, that indices of the SAM and ENSO would have skill in describing surface temperature anomalies, but it found mixed results. For example, the Southern Oscillation index (SOI) does explain SST and sea ice anomalies well over the Southern Ocean, but it does not have good skill at describing temperature anomalies on the Antarctic continent. Also, as Comiso (2000) and Shuman and Comiso (2002) discuss, the T_{IR} dataset is biased by the absence of data for days with cloud cover. Shuman and Comiso was the first study to directly compare T_{IR} and T_B data and generally found good agreement, but it only made comparisons at a few isolated locations with weather stations. Given the sparse distribution of Antarctic weather stations, it is desirable to further examine Antarctic climate with these satellite data. Other gridded products, such as the National Centers for Environmental Prediction–National Center for Atmospheric Research (NCEP–NCAR) reanalysis data, are significantly less reliable for Antarctic climate studies, especially for surface conditions (Marshall 2002b; Hines et al. 2000).

In this paper, we analyze the T_{IR} and T_B data in order to reduce uncertainties in interpreting either satellite dataset alone. First, we evaluate the T_{IR} data with methods that optimize the amount of variance that can be explained, in parallel to SS. Second, we use the two types of data to estimate the magnitude of microwave emissivity fluctuations. Next, the data are evaluated jointly using maximum covariance analysis (von Storch and Zweirs 1999). The results of this analysis increase confidence in the interpretation of both datasets in terms of surface temperature variability. We examine the relationship between surface temperature variability and atmospheric circulation through comparison of the empirical modes of the satellite datasets with NCEP–NCAR geopotential height data. We conclude that, overall, the SAM explains the greatest variance in Antarctic temperatures. However, the second most important influence is not simply described by ENSO, but rather, reflects a combination of patterns previously referred to as the Pacific–South American and wavenumber-3 patterns. We also suggest that blocking may be responsible for driving strong temperature trends in a little-studied region (0°–90°E) of East Antarctica.

2. Data

Surface temperature (T_{IR}) fields, at monthly resolution from January 1982 to December 1999, were derived for the Antarctic continent from thermal infrared channels of the Advanced Very High Resolution Radiometer (AVHRR) satellite as originally discussed by Comiso (2000). Comparison with available ground-based observations shows that T_{IR} data provide good estimates of the near-surface air temperature (T_a), although they may be cooler than the actual T_a under strong surface inversion conditions (Comiso 2000). In addition, monthly means of T_{IR} data have a clear-sky bias because infrared surface temperature estimates cannot be made in cloudy conditions. Since the net effect of clouds on surface temperature in the Antarctic is warming (e.g., King and Turner 1997), monthly cloud-free averages from the infrared observations tend to be cooler than in situ station observations by ~0.5 K (Comiso 2000). Originally constructed on a 6.25 km × 6.25 km polar stereographic grid, the T_{IR} data are averaged to a 25 km × 25 km grid so that they are coregistered with the passive microwave data. Anomalies are computed by subtracting the monthly climatology at each grid point.

Passive microwave brightness temperature (T_B) data used in this study are from the 37-GHz vertically polarized channel on the Scanning Multichannel Microwave Radiometer (SMMR) and Special Sensor Microwave Imager (SSM/I) satellite instruments from the same time period as the T_{IR} data. An important advantage of T_B data over T_{IR} data is that they can be obtained in all weather conditions. As discussed by SS, T_B data cannot be interpreted as a pure surface temperature signal, because the variability in microwave emissivity is not known. Attenuation of surface temperature changes through the penetration depth of the microwave emission—typically a few centimeters to a meter—means that amplitude is generally smaller for T_B variations than T_a variations. Surdyk (2002) emphasizes that changes in the snow temperature over the penetration depth have a much stronger influence on T_B variability than do emissivity changes. Surface melting, because of enhanced absorption of microwaves by liquid water during the melting event, and enhanced scattering after the snow refreezes (Zwally and Fiegles 1994), accounts for the largest emissivity-forced component of T_B. This was evaluated by SS and found to be important primarily along the coast and on the ice shelves. A melt-masked, monthly anomaly dataset of T_B on the Antarctic continent is derived in the same manner as described by SS except where noted below.

Previous site-specific studies have found good agreement between T_{IR}, T_B, and T_a measurements made in situ by automatic weather stations (AWS). On monthly time scales, differences among the data have been found to be less than 1 K at most locations, if T_B data are corrected for emissivity (Shuman and Comiso 2002). However, these reported differences do not take into account the clear-sky bias, as AWS temperatures were compared to T_{IR} data only on days when both observations were available. Shuman and Comiso (2002) also found evidence for consistent offsets between T_a and T_{IR} data, notably an ~4 K difference across all temperatures at the South Pole. This offset, which may be due to covering the South Pole at scan angles off of the nadir, is not significant in our analysis because the mean T_{IR} values are subtracted from each grid point. Also, the

T_B data are limited to areas north of 85°S, so the datasets are not directly compared in the South Pole region.

To examine the connections between the variability in Antarctic T_{IR} and T_B data and larger-scale SH atmospheric circulation variability, 1982–99 500-hPa geopotential height anomalies ($Z500$) poleward of 20°S, on a 2.5° × 2.5° latitude–longitude grid, are used from the NCEP–NCAR reanalysis (NRA; Kalnay et al. 1996; Kistler et al. 2001). Various biases, most importantly spurious multiannual trends, have been reported in these data (Hines et al. 2000; Marshall 2002b) but should be of little consequence to the purposes of this study. The 500-hPa level is the lowest standard pressure surface entirely above the surface of the ice sheet and the stable inversion layer, and NRA 500-hPa data compare more favorably to Antarctic station observations than do 850-hPa geopotential height or surface pressure fields (Marshall 2002b; Hines et al. 2000).

3. Methods

In section 4, principal component analysis (PCA) of T_{IR} anomaly data is performed. Empirical orthogonal function (EOF, spatial) patterns and principal components (PC, temporal variations) are computed for data from all months of the year and broken down by season, December–January–February (DJF), March–April–May (MAM), June–July–August (JJA), and September–October–November (SON). In section 5, the datasets are compared qualitatively, and in section 6, spatially and temporally varying microwave emissivity (ε) is estimated using the Rayleigh–Jeans approximation. Comparison of the T_{IR} and T_B fields through maximum covariance analysis (MCA) is used in section 7 to diagnose the common spatial–temporal signals in the two datasets. Heterogeneous regression maps are shown to illustrate the spatial patterns of the MCA modes, while expansion coefficients show temporal variations of the modes. Next, the leading T_{IR} and T_B expansion coefficients are compared through spectral analysis. In section 8, PCA is used to determine the leading patterns of variability in the atmospheric circulation at 500 hPa. Finally, regression analysis is performed among the various fields to show associated patterns.

4. The T_{IR} data

a. PCA of T_{IR} data

Applying PCA to the covariance matrix of monthly T_{IR} anomalies covering the Antarctic continent results in two modes with distinct eigenvalues that meet the separation criteria of North et al. (1982). The leading mode explains 52% of the variance in T_{IR}, while the second mode accounts for 9% of the variance. The first EOF, shown in Fig. 1a as a regression of T_{IR} anomaly data onto the first normalized principal component (T_{IR}–PC1; Fig. 1b) is associated most strongly with the high plateau of East Antarctica. Locally, high correlations in East Antarctica indicate that up to 80% of the variance in T_{IR} can be explained by this first mode, as determined by r^2 values. More moderate correlation of the same sign occurs over West Antarctica. Moderate correlation of opposite sign occurs on the northern reaches of the Antarctic Peninsula.

The second EOF (Fig. 1c) is centered on the Ross Ice Shelf and on the Marie Byrd Land region of the continent, where 40%–60% of the T_{IR} variance is explained. Most of West Antarctica is of the same sign, but the pattern changes sign over the Ronne–Filchner ice shelf (at 60°W) and most of East Antarctica. Some coastal areas near 120°E have the same sign as West Antarctica. Only a small fraction of the variance in East Antarctic temperatures can be explained by mode 2.

b. Seasonality of T_{IR} modes

While the leading patterns of tropospheric circulation variability, including the SAM, exist year-round in the SH middle and high latitudes (Cai and Watterson 2002; Gong and Wang 1999; Thompson and Wallace 2000), the range of Antarctic temperature variability is much larger in winter months than summer months (King and Turner 1997; Shuman and Stearns 2001). In T_{IR}, the standard deviation of July monthly means is about twice that of January monthly means, averaged over the continent. In winter, longwave radiation terms dominate the surface energy budget and strong surface inversions develop during clear and calm weather. Therefore, the surface temperature in winter is very sensitive to factors that disturb the inversion, particularly changes in cloudiness and winds (van den Broeke 2000b; Warren 1996). Since PCA modes are designed to maximize variance explained, the leading T_{IR} modes may be more characteristic of winter than of summer temperature variability.

PCA is performed on seasonal subsets of the data: austral summer (DJF), autumn (MAM), winter (JJA), and spring (SON). The short time series diminish the statistical significance of the modes compared to the full dataset; however, the following results can be supported. Mode 1 in every season dominates explained variance compared to the subsequent modes (Table 1), but explains slightly more variance in the transition seasons than the solstitial seasons. Compared to the full dataset, the mode-1 EOF pattern is most similar in spring and summer rather than the winter, as might be expected due to the larger variance in winter months. In spring and summer, a large part of the variance is explained over a broad region of East Antarctica, and most of the continent has the same sign. In autumn and winter, EOF 1 has considerably less coherent large-scale spatial structure. However, for all seasons, the correlation of the leading PC and the seasonal SAM index is about 0.5–0.6, as shown in Table 1, but is higher in the solstitial seasons. On the balance, these results suggest that

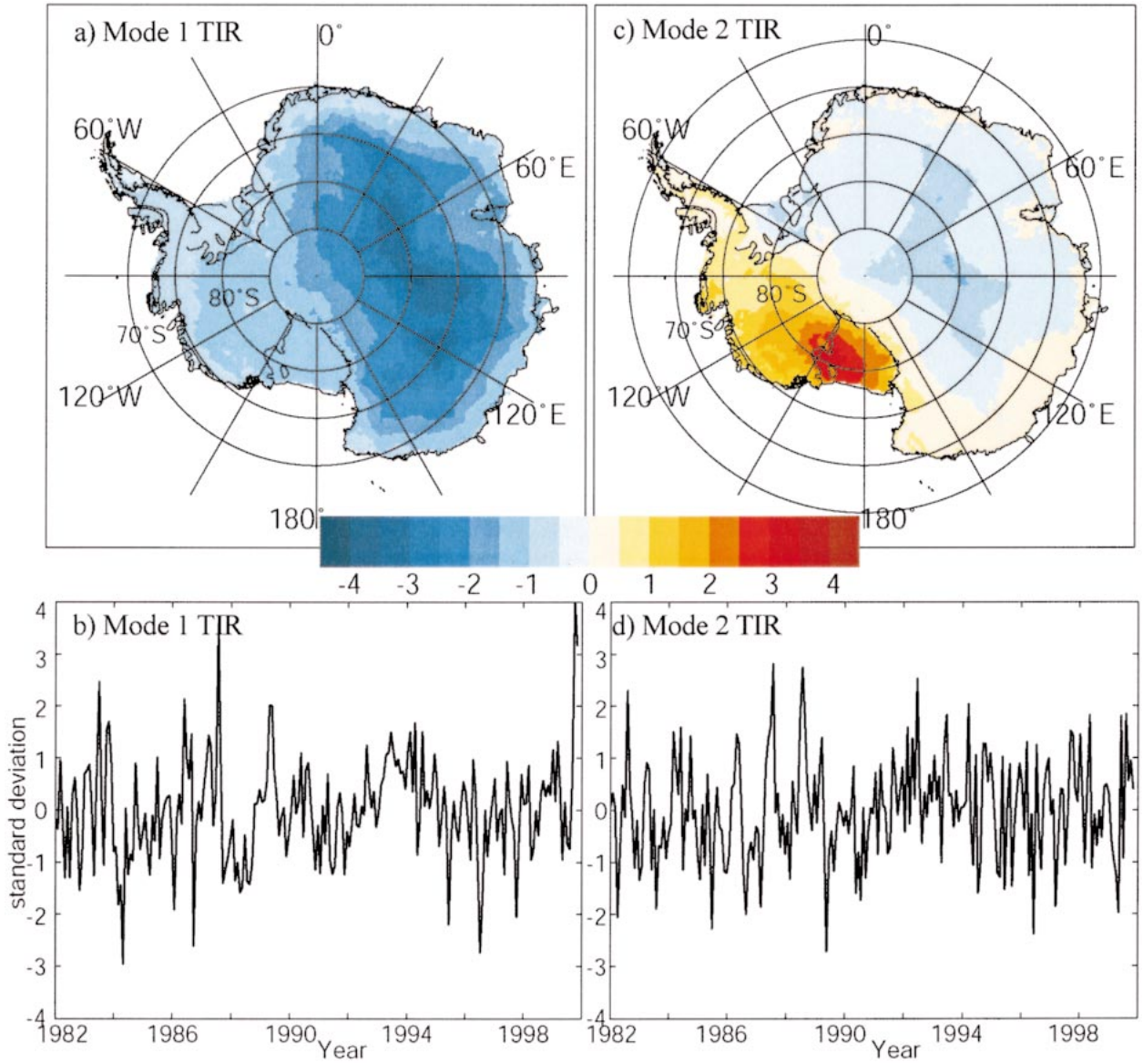


FIG. 1. Results from PCA of monthly 1982–99 T_{IR} anomaly data. (a) EOF-1 shown as a regression coefficient between each grid point and the first normalized PC. (b) The first normalized PC corresponding to the EOF pattern above. (c) As in (a) but for the second EOF. (d) As in (b) but for the second PC. Color scale is in K, corresponding to a typical anomaly associated with each mode, that is, the value of one positive standard deviation of the respective PC.

TABLE 1. Variance explained by first two PCA modes of T_{IR} and correlation with the SAM index by season.

Season	Mode 1 (%)	Mode 2 (%)	r(SAM index*, T_{IR}–PC1)
DJF	49	14	0.63
MAM	50	9	0.49
JJA	49	13	0.68
SON	57	7	0.50
Full year	52	9	0.57

* The SAM index is the first principle component of 500-hPa geopotential height as discussed in the text.

the year-round leading modes in T_{IR} are not heavily biased toward a particular season.

In all except JJA, mode 2 explains the most variance in the 90°W to 180° region of West Antarctica. The overall spatial patterns are less stable than in mode 1, and it is difficult to relate them to the full-year pattern. However, the anomalies are generally consistent with a combination of the zonal wavenumber-3 and PSA patterns of atmospheric circulation variability, as discussed in section 8. Since the atmospheric circulation patterns do exist throughout the year, and the traditional three-month seasonal breakdown is not very representative of Antarctica's seasons, it is reasonable to base interpre-

tations on the full-year dataset. With a 6-month seasonal breakdown, SS found similar leading modes in T_B for winter half (April–September) and summer half (October–March) years.

5. Comparison of T_{IR} modes with leading modes in T_B

Qualitatively, the leading modes of T_{IR} and T_B (SS) are similar. Mode 1 is characterized by covariance of the same sign throughout most of the continent, with the exception of some portions of the Antarctic Peninsula. Mode 2 has a center of action over the Ross Ice Shelf and adjacent inland areas in Marie Byrd Land. However, the differences in PCA modes of T_B and T_{IR} illustrate some contrasting features of the two datasets. Mode 1 explains 52% of the variance in T_{IR} but only 25% of the variance in T_B. Mode 2 explains 9% of the variance in T_{IR} and 18% of the variance in T_B. This can be partially attributed to differences of spatial and temporal autocorrelation. While the monthly T_{IR} data have a typical (for an atmospheric variable) lag-1 autocorrelation coefficient of $\sim$0.27, T_B data are temporally autocorrelated at $\sim$0.60, resulting in less separation between the T_B modes. Spatially, T_B data are quite variable, due to regional differences in snow physical properties, while the T_{IR} data vary on much larger length scales, reflecting more closely the surface temperature.

6. Emissivity variability at microwave wavelengths

The emissivity parameterization helps to quantify the effects of the physical properties of the snow layer from which the T_B signal emanates. Here T_B is the physical temperature of the snow times its emissivity (ε), integrated over the penetration depth (Zwally 1977; Surdyk 2002). This is the Rayleigh–Jeans approximation, allowing a calculation of 37 GHz ε through the relation $\varepsilon = T_B \times T_{IR}^{-1}$. On annual mean or longer time scales, using T_{IR} data in place of the physical temperature of the snow is valid, as mean emissivity changes little and the mean annual surface temperature approximates the mean physical temperature of the snow (Zwally 1977; Surdyk 2002). Coregistration of T_{IR} and T_B data enables a map of ε to be calculated based on the 1982–99 means (Fig. 2a). The pattern shows a spatial variation in ε of 0.25. This includes the melt areas, which have not been excluded from the T_B data in this section of our study. Although the average spatial variations in ε were removed from the analysis of SS by the use of anomalies and the masking of melt zones, the influences of spatially varying ε do affect the appearance of EOF–regression maps in some areas, especially near the margins of the ice sheet. Melting can temporarily make the magnitude of T_B anomalies greater than that of T_{IR} anomalies due to the high absorption of the liquid water.

Temporally, ε is negatively correlated with the annual cycles of T_{IR} and T_B when averaged over the continent,

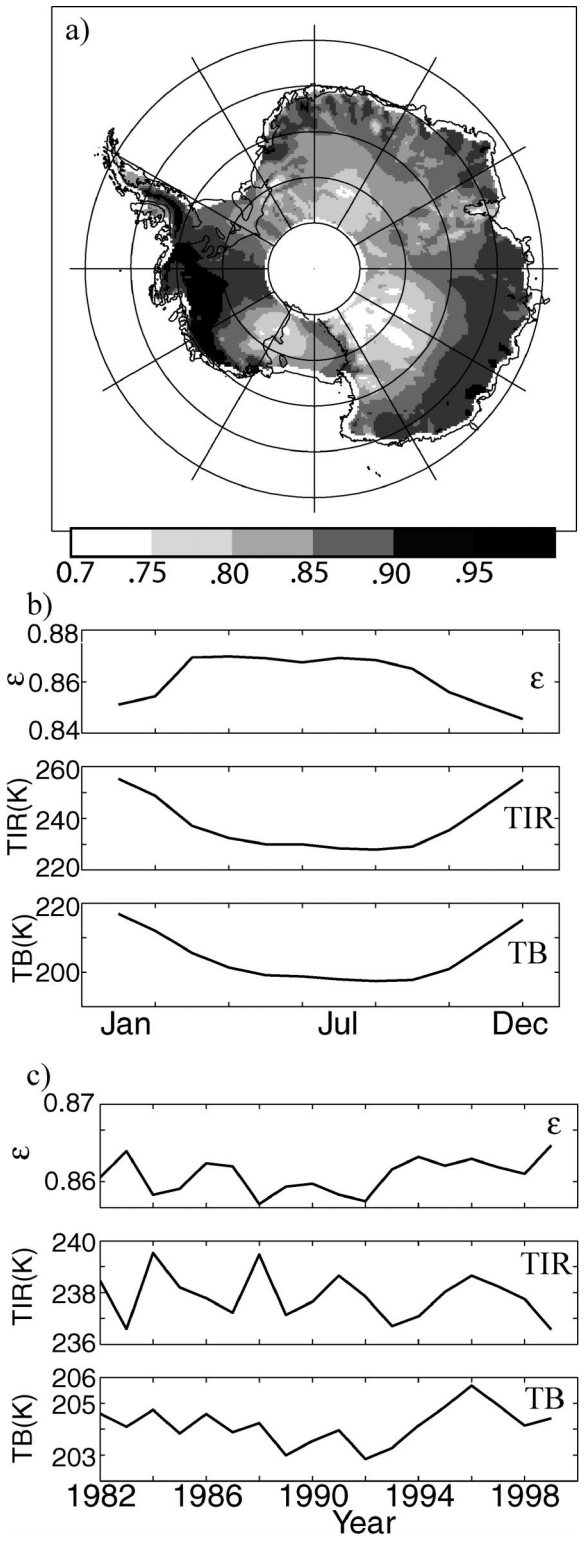


FIG. 2. (a) Mean microwave emissivity at 37-GHz vertical polarization of the Antarctic ice sheet based on 1982–99 mean values of T_{IR} and T_B. (b) Continent-averaged annual cycles of emissivity, T_B, and T_{IR} based on 18-yr means of each month. (c) Continent-averaged interannual variations of emissivity, T_B, and T_{IR}.

TABLE 2. Correlation coefficients among PCs of T_B and T_{IR} considered separately and expansion coefficients from MCA. Sign is ignored.

	T_B–PC1	T_{IR}–PC1	T_B–PC2	T_{IR}–PC2	T_B–EP1	T_{IR}–EP1	T_B–EP2	T_{IR}–EP2	T_B–EP3	T_{IR}–EP3
T_B–PC1	1	0.46	0	0.11	0.81	0.47	0.21	0.05	0.01	0.04
T_{IR}–PC1		1	0.05	0	0.68	0.99	0.01	0.11	0.00	0.02
T_B–PC2			1	0.48	0.03	0.04	0.96	0.56	0.14	0.02
T_{IR}–PC2				1	0.09	0.02	0.54	0.96	0.13	0.36
$T_B \times 1$					1	0.70	0.07	0.00	0.07	0.00
$T_{IR} \times 1$						1	0.00	0.09	0.00	0.00
$T_B \times 2$							1	0.61	0.11	0.00
$T_{IR} \times 2$								1	0.00	0.15
$T_B \times 3$									1	0.78
$T_{IR} \times 3$										1

and the apparent magnitude of the seasonal change in ε is about 0.02 (Fig. 2b). However, this magnitude is partly an artifact, as the T_B annual amplitude (19 K) is 30% less than the T_{IR} annual amplitude (27 K). This ~30% attenuation is indicative of an average penetration depth of less than 1 m (Surdyk 2002). Thus, the true seasonal variation in ε must be less than the apparent magnitude obtained when the attenuation is ignored (Surdyk 2002). An attenuation map can be produced by dividing the average magnitude of the T_B annual cycle by the magnitude of the T_{IR} annual cycle (not shown). By inference, the more damped the annual cycle, the deeper the penetration depth. The spatial pattern in attenuation is highly correlated with the emissivity map, showing that low emissivity corresponds to shallow penetration depth and vice versa, consistent with the theory that grain size is the dominant factor affecting both parameters (Surdyk 2002). Annual mean time series of T_{IR} and T_B are positively correlated, while the emissivity is anticorrelated with T_B and T_{IR} (Fig. 2c). The calculated interannual range in the value of ε is on the order of 0.01 with a standard deviation of 0.0034. This value of 0.01 is likely close to the real range in ε over the penetration depth because the annual mean surface temperature (from T_{IR}) approximates the annual mean temperature at depth. Thus, given the mean annual T_{IR} value of 239 K and the mean ε of 0.86, the interannual standard deviation in ε accounts for only ~0.8 K of (microwave brightness) temperature difference, well within the uncertainties of both datasets (Shuman and Comiso 2002) and well below the magnitude of temperature anomalies that are explained by our modes in section 7.

7. Maximum covariance analysis of the datasets

Maximum covariance analysis (MCA) optimizes the covariance explained by pairs of structures in two datasets. Bretherton et al. (1992) and Wallace et al. (1992) provide a detailed discussion of the methodology, which is adhered to below. The name singular value decomposition (SVD) is often applied to the entire method; here it is only used in reference to the algorithm used in extracting empirical structures via cross-covariance matrix decomposition.

First, the cross-covariance matrix of T_{IR} and T_B anomaly fields is computed (with melting pixels masked as in SS). The expansion coefficients of the T_{IR} and T_B fields are found by projecting the singular vectors from SVD onto the original gridpoint data of the respective field. These expansion coefficients are then normalized, and either field can be regressed upon them to display spatial structure. In the case that the T_{IR} (T_B) field is regressed upon a T_{IR} (T_B) expansion coefficient, the map is known as a homogeneous regression map and, in the case that the T_{IR} (T_B) field is regressed upon a T_B (T_{IR}) expansion coefficient, the map is referred to as a heterogeneous regression map.

No formal method has been developed for determining the significance of MCA modes, but some tests apply. The squared covariance fraction (SCF) of each mode is an indication of the fit between the two datasets. Another indication is the correlation coefficient between each mode's pair of expansion coefficients. Additionally, the cross-covariance matrix can be tested for relatedness (before applying MCA) with root-mean-square covariance [RMSC, square root of the squared covariance (SC) after dividing by the product of the variance of the two datasets]. The high T_{IR}–T_B RMSC of 0.22 implies strongly coupled fields that are suitable for MCA, as RMSC of 0.1 or greater is a typical guideline for strong correlation (Wallace et al. 1992).

MCA applied to T_{IR} and T_B fields yields three significant modes, with a SCF of 77%, 11%, and 5%, respectively. The modes' three pairs of expansion coefficients (T_{IR}–EP1 and T_B–EP1; T_{IR}–EP2 and T_B–EP2; T_{IR}–EP3 and T_B–EP3) are correlated at 0.70, 0.61, and 0.78, respectively (Table 2). In both T_{IR} and T_B, the set of homogeneous regression maps for the leading two modes are almost identical to the leading spatial patterns from PCA of the datasets considered separately.

The set of heterogeneous regression maps, displayed in Fig. 3, are similar to cross-regressions and show the anomaly in the field on the map associated with one standard deviation of the opposite field's expansion coefficient. For mode 1, the T_{IR} field regressed onto T_B–EP1 (Fig. 3a) is similar to the first EOF of T_{IR} anomalies, although less variance in T_{IR} is explained (23%). The T_B field regressed onto T_{IR}–EP1 (Fig. 3b) produces a pattern that is more smoothly varying than the first EOF

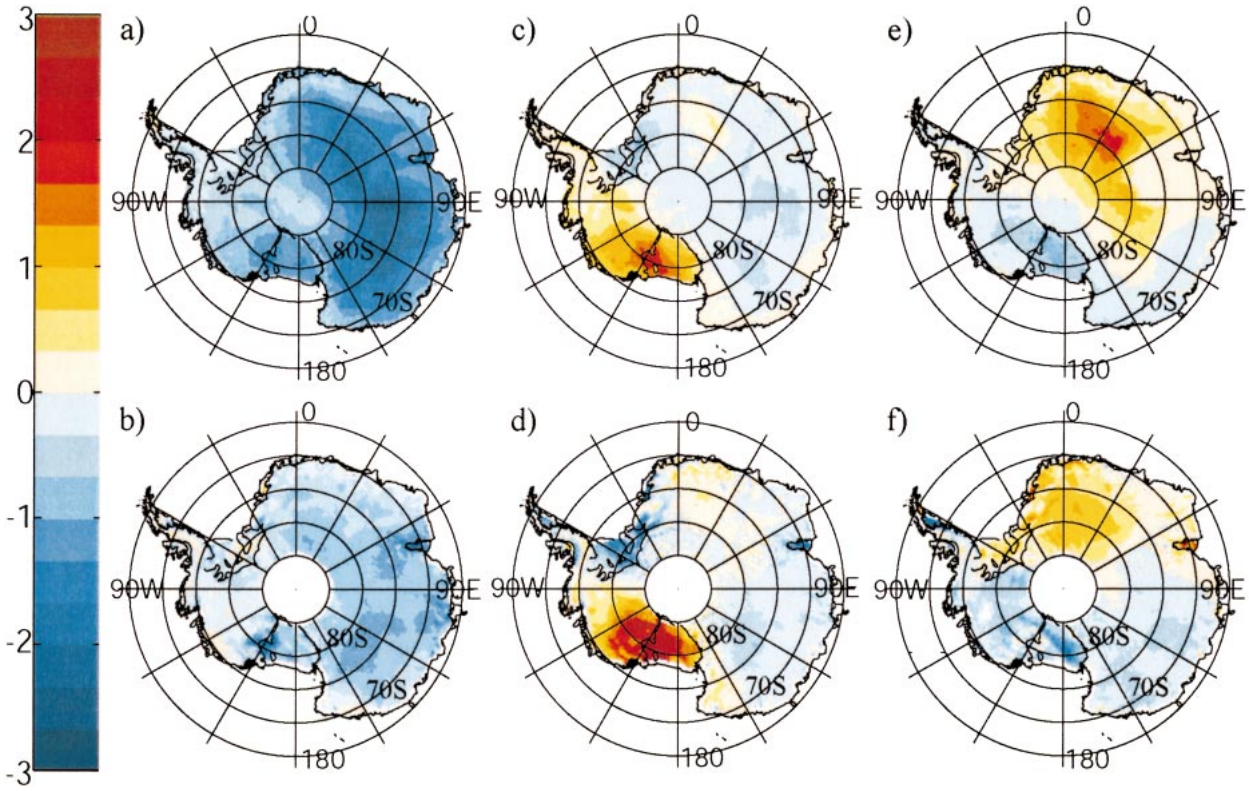


FIG. 3. Heterogeneous regression maps from MCA of T_{IR} and T_B fields. (a), (c), (e) Covariances from the T_{IR} field regressed upon the first, second, and third normalized T_B expansion coefficients, respectively; (b), (d), (f) covariances from the T_B field regressed upon the first, second, and third normalized T_{IR} expansion coefficients, respectively. Color is in units of K, corresponding to one standard deviation of the respective expansion coefficient.

of T_B (see SS02, Fig. 2a), and explains 11% of the variance. The amplitudes in Fig. 3b are smaller than those in Fig. 3a, implying attenuation of the surface temperature signal through the penetration depth. Also, where the covariance in Fig. 3b is the greatest in East Antarctica, notably between 75° and 80°S, the emissivity is lowest, as shown in Fig. 2a, consistent with shallow microwave penetration depths and little attenuation. In Fig. 3c, the T_{IR} field is regressed onto T_B–EP2, explaining 3% of the variance in T_{IR}, and the resulting pattern is similar to the second EOF of T_{IR}. Likewise, regression of the T_B field onto T_{IR}–EP2 (Fig. 3d), explains 4% of the variance, and produces a heterogeneous map similar to the second EOF of T_B. In this pair of maps, the amplitudes are comparable in magnitude, consistent with little attenuation, a shallow penetration depth, and low emissivity in the Ross Sea sector of Antarctica (Fig. 2a).

A third mode is diagnosed with MCA that was not prominent in the PCA results for the datasets when considered separately (although this third MCA mode correlates well with the fourth mode in T_{IR} data alone and the fifth mode in T_B data alone). It is retained for discussion because it projects onto the linear trends in the T_{IR} and T_B datasets and is reproducible, as discussed below. Because the T_{IR}–EP3 and T_B–EP3 time series have upward trends (see Fig. 4c below), these time series and the T_{IR} and T_B gridpoint data are both detrended prior to the construction of the heterogeneous regression maps in order to avoid spurious correlations. If the trends are retained in the time series, the heterogeneous maps of mode 3 look very similar to annual mean trends in the T_{IR} and T_B datasets [see Kwok and Comiso (2002), Fig. 2 for trends in T_{IR}]. Therefore, care must be taken not to include spurious correlations of unrelated trends in the maps. The map with T_{IR} data regressed onto T_B–EP3 explains 4% of T_{IR} variance and shows positive anomalies in T_{IR} throughout much of East Antarctica (Fig. 3e). The anomalies of greatest magnitude occur from 0° to 60°E. Similarly, the map of T_B data regressed onto T_{IR}–EP3, explains 4% of T_B variance and shows positive T_B anomalies, but of weaker magnitude, in the same area of East Antarctica.

The results outlined above strongly imply that there is meaningful covariance between T_{IR} and T_B datasets. However, it must be established with confidence that the correlations have not arisen by chance. As a first test of reproducibility, T_{IR} and T_B anomaly data are detrended prior to MCA. In this case, the same three modes are produced, but without the trend in the third mode. Second, as a test of the statistical robustness of the MCA results, the T_{IR} and T_B datasets are divided into subsets.

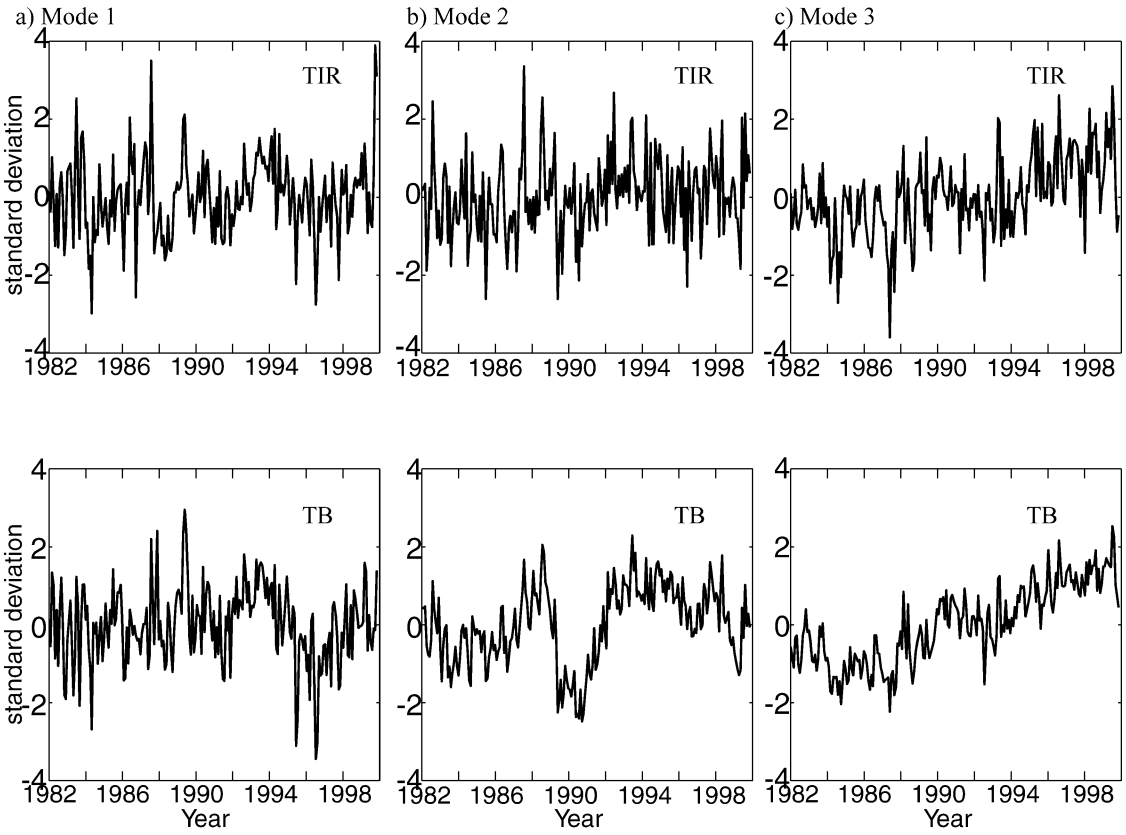


FIG. 4. Expansion coefficients of the first three MCA modes [(a)–(c), respectively] corresponding to the heterogeneous maps in Fig. 3.

Data for odd months only are used, and then data for even months only. Odd and even month RMSC, SC, SCF, and correlation coefficients between expansion coefficient pairs are comparable in magnitude to the statistics for the full datasets for each of the first three modes (Table 3), indicating that the first three modes meet reproducibility criteria.

Another test of statistical robustness, based on the following Monte Carlo procedure, further demonstrates the strong relationships between T_{IR} and T_B anomaly data. Following the method of Wallace et al. (1992), the temporal order of the T_{IR} field is scrambled randomly while the order of the T_B field remains unchanged. RMSC, SC, SCF, and correlation coefficients are computed for each of 1000 random runs (Table 3). The significance of the observed runs and the subsets clearly stands out above the random runs. The observed squared covariance is an order of magnitude larger than the mean SC of the scrambled runs, the SCF of the observed first mode is about two standard deviations above the mean of the scrambled runs, and the correlation coefficients among expansion coefficient pairs for the first three modes are well above the mean of the scrambled runs. Finally, the RMSC values of the random datasets are smaller than the value of 0.1 that would indicate strong correlation. This leaves little doubt as to the significance, above the 99% confidence level, of the leading MCA modes and the strong relationship between the T_{IR} and T_B fields.

TABLE 3. Summary of MCA reproducibility tests and Monte Carlo results.

Run	SC	RMSC	SCF (%) modes (1, 2, 3)	r (T_Bx, T_{IR}x) modes (1, 2, 3)
		1000 Scrambled runs		
Highest	6.3×10^4	0.098	(82, 31, 18)	(0.34, 0.39, 0.38)
Lowest	1.6×10^4	0.050	(29, 04, 03)	(0.14, 0.16, 0.16)
Mean	3.1×10^4	0.068	(55, 14, 09)	(0.23, 0.28, 0.28)
Std dev	7.2×10^3	0.008	(10, 04, 02)	(0.03, 0.04, 0.04)
		Observations		
	3.28×10^5	0.223	(77, 11, 05)	(0.70, 0.61, 0.78)
		Odd months only		
	3.27×10^5	0.241	(77, 11, 05)	(0.74, 0.54, 0.80)
		Even months only		
	3.32×10^5	0.222	(70, 16, 06)	(0.69, 0.67, 0.75)

The three pairs of normalized expansion coefficients show the time variability of the three surface temperature modes from MCA (Fig. 4). The expansion coefficients of the first two modes are well correlated with the original PCs from each dataset considered separately, as can be seen in Table 2. For instance, T_{IR}–PC2

correlates with T_{IR}–EP2 at $r = 0.96$ and similarly high correlations exist for the other matches. Therefore, the time series that explain the most covariance between the T_{IR} and T_B datasets also explain the most variance in the individual datasets.

The power spectra of each expansion coefficient time series are estimated with a Hanning window providing 13 degrees of freedom (Fig. 5). Also shown in Fig. 5 is the theoretical spectrum, with 95% confidence limits, for the first-order red noise autoregressive process [AR(1)] that provides the best fit to each time series (von Storch and Zweirs 1999). There are no significant spectral peaks at the 99% confidence level. Each time series is consistent with red noise, but with different degrees of ''redness,'' which can be quantified by comparing the AR(1) coefficients from the best-fit theoretical spectrum.

If it is assumed a priori that T_{IR} spectra provide a direct measure of the variability in the surface temperature, then T_B spectra should reflect the attenuation of that variability at depth (see Fig. 10 of Surdyk 2002). Because both the penetration depth (as discussed in section 6) and the surface forcing (as indicated by the T_{IR} and *Z*500 modes) are spatially variable, the degree of attenuation and ''memory'' in the T_B time series will also be spatially variable, and this is reflected in differences among the pairs of AR(1) coefficients. The MCA procedure is therefore unable to completely remove the effect of spatially varying snow structure, but it does show that the T_B and T_{IR} have a common forcing—the surface temperature. The strong covariance between the datasets suggests that effects of the clear sky bias in T_{IR} data and snow emissivity influences on T_B data do not mask the underlying large-scale modes of surface temperature variability on monthly time scales or longer.

8. Influence of atmospheric circulation on Antarctic temperatures

a. Principal component analysis of 500-hPa geopotential height anomalies

To define the leading patterns of SH atmospheric circulation during the time period of this study in a consistent manner with the analysis of temperatures in section 4, PCA is applied to monthly *Z*500 anomaly data poleward of 20°S. For equal-area weighting, the data are weighted by the square root of the cosine of their latitude prior to analysis. The original unweighted *Z*500 data are regressed against each normalized PC, showing anomalies corresponding to one standard deviation of the corresponding PC. Three patterns of interest are resolved, explaining 24%, 12%, and 10% of the (weighted) variance, respectively. The latter two patterns are not well separated under the criteria of North et al. (1982). However, they have been reported by a number of studies and found in many different datasets (Carleton 2003; Cai and Watterson 2002; Mo and White 1985). Our results are consistent with the definition of the first *Z*500 pattern as the SAM (Fig. 6a), the second *Z*500 pattern as the Pacific–South American (PSA) pattern (Fig. 6b), and the third pattern as the zonal wavenumber-3 pattern (Fig. 6c) as named by other studies (Cai and Watterson 2002; Mo 2000). The third pattern is sometimes called the PSA-2 pattern (Mo 2000). The signs are displayed for consistency with the modes in T_{IR} and T_B and the regression patterns discussed below.

b. Correlations and regression patterns

Associations between Antarctic temperatures and patterns of atmospheric circulation can be illustrated with a variety of methods. The temporal correlations among the PCs of *Z*500, the expansion coefficient pairs from MCA, and the SOI are summarized in Table 4. The first PC of *Z*500 forms a representative SAM index (Thompson and Solomon 2002) and it has strong correlation with the first MCA mode [$r(T_{IR}$–EP1, SAM) = 0.58 and $r(T_B$–EP1, SAM) = 0.61] and weak correlation with the other modes. The *Z*500 PC2 correlates well with the SOI ($r = 0.43$) and has moderate correlation with both modes 1 and 2 in T_{IR} and T_B. The *Z*500 PC3 has a weak correlation with the SOI and the best correlation with mode 2 in T_{IR} and T_B [$r(T_{IR}$–EP2, *Z*500 PC2) = 0.39 and $r(T_B$–EP2, *Z*500 PC2) = 0.32]. Mode 3 from MCA has only weak correlation with the *Z*500 patterns.

Regression patterns reinforce the connections implied by the various correlation coefficients. Since T_{IR} and T_B data are highly correlated, regressions involving *Z*500 data discussed here are made only with T_{IR} data for illustration. Regressions of T_{IR} data onto the normalized PCs of the three leading *Z*500 patterns are shown in Fig. 7. The SAM explains 17% of the variance in T_{IR} anomalies (Fig. 7a), the PSA pattern explains 6% of the variance (Fig. 7b), and the zonal wavenumber-3 pattern explains 3% of the variance (Fig. 7c). The first regression pattern is quite similar to the first T_{IR} EOF (Fig. 1a). During the positive phase of the SAM, relatively strong westerlies encircle Antarctica near 60°S, which tends to enhance warm air advection over the northern peninsula, while the cool anomalies on the rest of the continent are indicative of adiabatic cooling (Thompson and Wallace 2000).

As seen in Fig. 7b, in East Antarctica, the PSA pattern explains much less variance in surface temperature than does the SAM, but the spatial structures of temperature anomalies are generally similar. In the peninsula and most of West Antarctica, the PSA pattern explains variability of the same sign as in East Antarctica. However, the PSA pattern is associated with temperature anomalies of opposite sign near 75°S, 150°W, consistent with the anticyclonic 500-hPa height anomaly centered near 60°S, 125°W (Fig. 6b).

The PSA regression pattern resembles the regression of T_{IR} data upon the SOI that was shown in Fig. 3b of

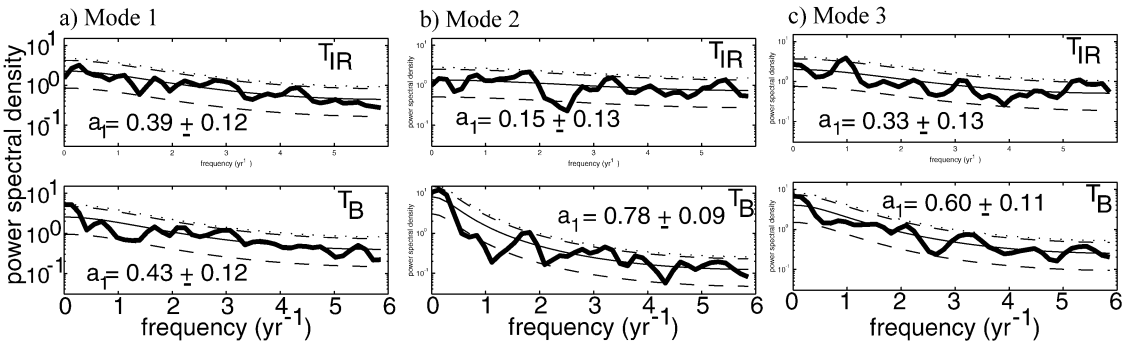


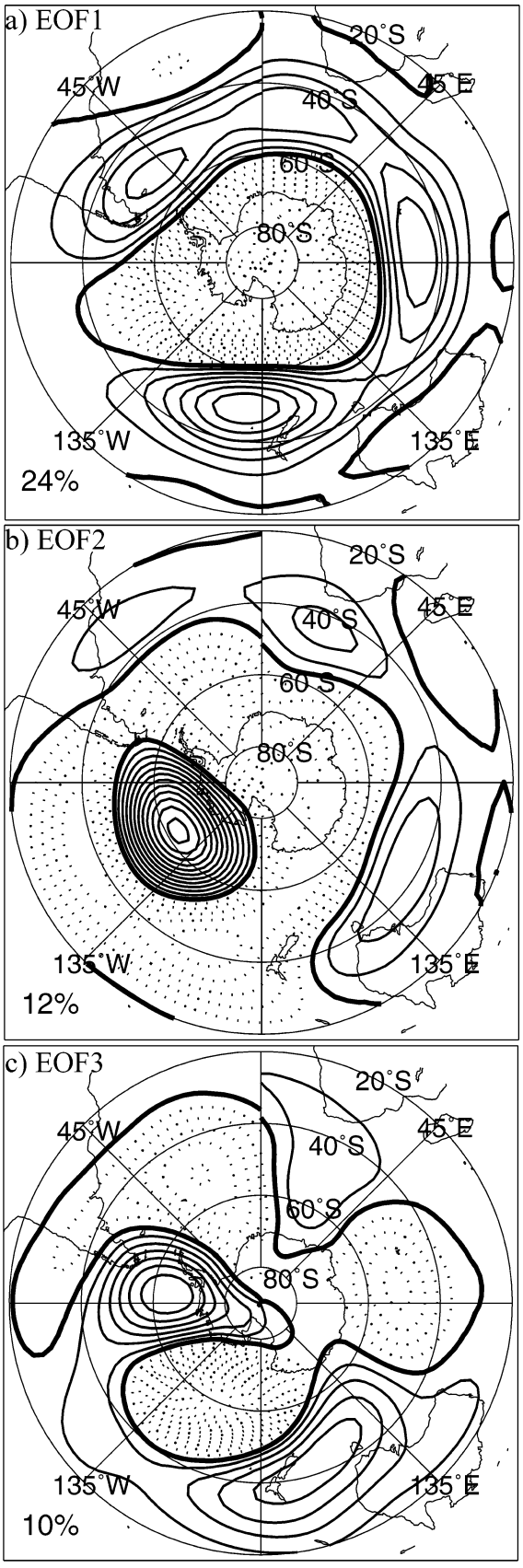
FIG. 5. Power spectra (heavy solid lines) of the MCA expansion coefficients shown in Fig. 4. for (a) first, (b) second, and (c) third modes. Also shown (thin solid lines) are the theoretical spectra of the AR(1) (red noise) process that best fits the original time series. The 5% and 95% confidence levels are indicated by dashed lines; 99% confidence level by dotted lines. The coefficients a_1 of the best-fit spectra are also indicated inside the plots.

Kwok and Comiso (2002). For comparison to the PSA pattern in geopotential height, we show the pattern in *Z*500 associated with the Southern Oscillation, formed by regressing *Z*500 data onto the negative SOI to illustrate anomalies typical of El Niño years (Fig. 8). The zonally elongate north–south dipole structure over the far southern Pacific closely resembles the ENSO warm minus cold year mean 500-hPa height difference of Renwick and Revell (1999, Fig. 5) and the ENSO-associated patterns of other studies (e.g., Mo and Higgins 1998; Kidson 1999; Bromwich et al. 2003). In contrast to the 6% of T_{IR} variance explained by the PSA pattern, our calculations indicate that only 0.5% of the variance in T_{IR} anomalies on the Antarctic continent is explained by the regression of T_{IR} data upon the SOI. However, circulation indices such as the SOI may not adequately capture the variability that is truly associated with ENSO (Carleton 2003; Kidson and Renwick 2002). ENSO-related forcing in the Tropics is thought to project primarily onto the PSA pattern of variability on interannual to interdecadal time scales in the middle to high latitude SH geopotential height field (Cai and Watterson 2002; Garreaud and Battisti 1999).

It is interesting that the pattern in Antarctic T_{IR} explained by the PSA pattern also resembles the correlation (although of opposite sign) of winter temperatures at Faraday Station with the T_{IR} gridpoint data [King and Comiso (2003) Fig. 1]. Although King and Comiso suggested that the climate variability of the Antarctic Peninsula is unrelated to the rest of Antarctica, our results show a connection through the PSA pattern. As illustrated in Fig. 6b, the strong anticyclonic anomaly in the far southeastern Pacific is accompanied by low geopotential height over most of Antarctica. Thus there is a connection between the anomalous meridional advection along the western peninsula implied by the southeastern Pacific center of action and contemporaneous decreases in geopotential height over most of the continent.

As shown in Fig. 7c., the wavenumber-3 pattern explains positive temperature anomalies over the same area of West Antarctica associated with the second surface temperature mode, consistent with the strong positive and negative height anomalies centered along 60°S near 90°W and 155°W, respectively. The wavenumber-3 pattern also explains weak temperature anomalies of the same sign in East Antarctica, unlike the second temperature mode, in which West and East Antarctica are out of phase. This is likely due to the additional influence of the PSA pattern on the second surface temperature mode. While the PSA pattern in *Z*500 is associated with geopotential height fluctuations over East Antarctica, the wavenumber-3 pattern has very little correlation with heights over this region.

As a consistency check, *Z*500 data are now regressed upon T_{IR} expansion coefficients from the leading MCA modes. As above, the time series of mode 3 are detrended prior to regression. Although T_{IR} expansion coefficients are used here for illustration, regressions involving T_B expansion coefficients are very similar. The first *Z*500 regression pattern (Fig. 9a) closely resembles the SAM pattern in *Z*500 (Fig. 6a), especially in the Eastern Hemisphere. In the west, the PSA (Fig. 6b) pattern appears to have an influence on both the first and second regression patterns (Figs. 9a and 9b). However, the second regression pattern (Fig. 9b) most strongly correlates with the *Z*500 wavenumber-3 pattern (Fig. 6c). The third regression pattern (Fig. 9c) resembles the SAM, but has anomalies of much weaker magnitude. The positive anomalies in East Antarctica near 45°E are suggestive of a ridge in the midtroposphere extending inland through East Antarctica. This resembles the wintertime blocking episodes discussed by Hirasawa et al. (2000) and Enomoto et al. (1998), who documented warm, moist air being pumped from the north all the way to the polar plateau. The positive temperature anomalies associated with the third mode (Fig. 3e) correspond to the ''strip'' of annual mean warming on the East Antarctic ridge observed by Kwok and Comiso (2002). Inspection of monthly trends shows that in winter months, warming trends in T_{IR} occur over a much broader area than this strip, and Comiso (2000) de-



scribed several anomalously warm July episodes in East Antarctica. Since the winter months account for most of the trend in the mode 3 expansion coefficients, it is likely that the events explained by mode 3 and these warm episodes are part of the same phenomenon.

9. Discussion and conclusions

Previous studies have shown that Antarctic 37-GHz T_B data and T_{IR} data are both well correlated with surface air temperatures (Shuman and Comiso 2002; Surdyk 2002; Shuman and Stearns 2001; Comiso 2000), but this is the first to fully examine the consistency of these relationships across the continent. The PCA and MCA results demonstrate that the most important empirical modes in the T_{IR} and T_B datasets are well correlated with each other. The strength of the connection between T_{IR} and T_B increases confidence in the quality of both datasets. A general difference between the datasets is that the T_B data are more spatially varying and more temporally autocorrelated than the T_{IR} data because of the dependence of T_B on both emissivity and temperature integrated over a layer of snow and firn. The degree of dampening of T_B signals depends on the period of the surface forcing and the penetration depth of the microwave emission. Lower emissivity regions indicate a shallow penetration depth and relatively high amplitude T_B fluctuations, which results in the highest T_{IR} and T_B covariance. The spatial pattern in emissivity apparently changes very little in time, and most likely represents spatially differing snow and firn structures.

Spatial and temporal patterns of T_{IR} and T_B variability, and more generally, surface temperature variability, in Antarctica are consistent with well-documented patterns of variability in extratropical SH atmospheric circulation. It is clear that the most important influence on Antarctic temperature anomalies from month-to-month to interannual time scales is the SAM. This first mode is well separated from other modes in both *Z*500 data and the satellite datasets. Looking at data from all months, strong temperature trends associated with this mode are not seen. However, inspection of trends by month over the length of the record shows that T_{IR} observations are consistent with a late spring and summer cooling trend, possibly driven by an increasing tendency of the SAM to stay in its positive phase during these seasons (Thompson and Solomon 2002).

The PSA pattern has an influence on the first two surface temperature modes. The wavenumber-3 pattern of variability, however, has a relatively stronger influ-

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FIG. 6. The leading modes in monthly 500-hPa geopotential height, 1982–99. The (a) first, (b) second, and (c) third EOFs, shown as the *Z*500 data regressed upon the leading normalized PCs. Percentage of variance explained indicated at lower left. Contour interval 5 m, zero contour (heavy solid line), negative contours (dotted), positive contours (solid).

TABLE 4. Summary of correlation coefficients among modes in *Z*500, SOI, and MCA expansion coefficients. Sign is ignored.

	SOI	T_B–EP1	T_{IR}–EP1	T_B–EP2	T_{IR}–EP2	T_B–EP3	T_{IR}–EP3
*Z*500-PC1	0.16	0.61	0.58	0.06	0.05	0.06	0.15
*Z*500-PC2	0.43	0.25	0.31	0.24	0.17	0.06	0.06
*Z*500-PC3	0.18	0.09	0.16	0.32	0.39	0.02	0.00
SOI		0.02	0.04	0.19	0.14	0.06	0.11

ence on the second surface temperature mode, shown by its association with large temperature anomalies in the West Antarctic sector inland of the Ross and Amundsen Seas. Anomalies of opposite sign in East Antarctica suggest that the PSA pattern exerts a stronger influence there. Since ENSO-related variability projects primarily onto the PSA and the wavenumber-3 patterns, Antarctic climate records often show ENSO-like spectra (SS; Ichiyanagi et al. 2002; Bromwich and Rogers 2001; White et al. 1999).

Some persistent trends in the available satellite record are associated with the third mode, which cannot, within the scope of this study, be clearly linked to the principal patterns of atmospheric circulation variability. However, blocking events over inland East Antarctica have been documented with station data (Hirasawa et al. 2000; Enomoto et al. 1998) and provide a plausible explanation for the trends in temperature and the pattern seen in *Z*500 on our regression map (Fig. 9c). During these episodes, rises of ~40 K can occur in two days or less at remote interior stations such as Dome Fuji and Plateau, and T_a can take more than a month to return to its value before the rise (Enomoto et al. 1998; Kuhn et al. 1973). The upward trend in the T_{IR} and T_B expansion coefficients comes primarily from the winter months, when the blocking episodes most often occur and the surface temperature is extremely sensitive to circulation changes. In addition, changes in cloud cover and winds associated with blocking will destroy the surface inversion, adding to the magnitude of the surface temperature anomalies (Hirasawa et al. 2000). At present, however, the satellite record is too short to establish the long-term significance of the trends, and the monthly temporal resolution of this study limits our ability to further characterize the causes of variability in the third temperature mode.

The results of this study show that the surface temperature variability of Antarctica is well represented by both the T_{IR} and T_B datasets. Nonetheless, it is important to note that biases in the T_{IR} data associated with cloud cover, and in the T_B data associated with attenuation and possible emissivity changes, have not been completely removed. Ongoing improvements to the data include a technique for filling in cloud gaps in infrared observations with emissivity-corrected T_B observations (C. A. Shuman 2002, personal communication) and a method suggested by Winebrenner et al. (2004) that models T_B data on the basis of T_a variations and which could lead to an improved parameterization of snow properties and their influence on T_B. Analyses of updated datasets may result in minor changes to the empirical modes we have calculated.

Currently, a gap exists between our understanding of Antarctica's short instrumental and satellite records and deep ice cores from Antarctica (e.g., Petit et al. 1999; Morgan et al. 2002). Future work will include evaluating the stability of the temperature and circulation modes discussed here on longer time scales. Some prior work hints that these modes operated in the past. For instance, long ice core paleoclimate records from locations spread thousands of kilometers apart in East Antarctica are well correlated with each other, consistent with our first mode, while records from West Antarctica and some East Antarctic cores appear to reflect local climate, consistent with our second mode (Watanabe et al. 2003; Steig et al. 2000). East Antarctic isotopic records from

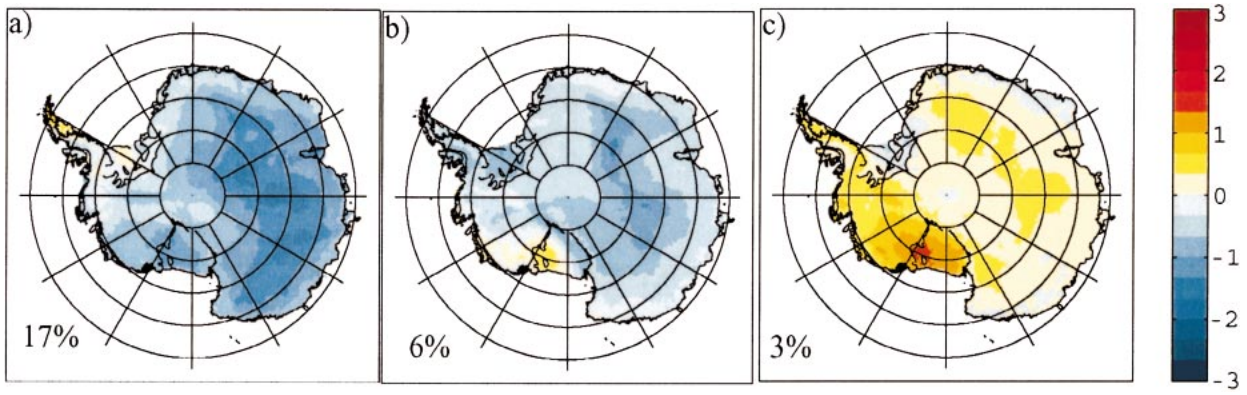

FIG. 7. Regressions of T_{IR} gridpoint data upon the first three *Z*500 normalized PCs [(a)–(c), respectively]. Percentage of variance explained indicated in lower left. Color scale is in K, corresponding to one standard deviation of the respective PC time series.

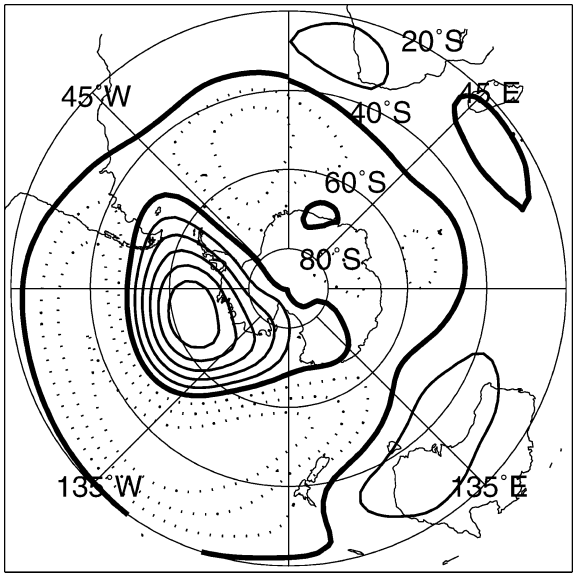

FIG. 8. Regression of 1982–99 monthly *Z*500 data upon the SOI with sign reversed to show anomalies typical of the ENSO warm phase. Units, contours as in Fig. 6.

the ice cores may be tied to the SAM over a large area (Noone and Simmonds 2002), while West Antarctic ice core records would be expected to be strongly linked to circulation variability in the Southern Pacific, which in turn is teleconnected to the tropical Pacific during strong El Niño and La Niña events (Bromwich et al. 2003). Century-scale reconstructions of the major modes of SH atmospheric circulation from tree rings (Jones and Widmann 2003) and a network of intermediate-depth Antarctic ice cores (Mayewski 2003) will help to fill in the gap between our understanding of modern climate variability and our theories of past climate variations.

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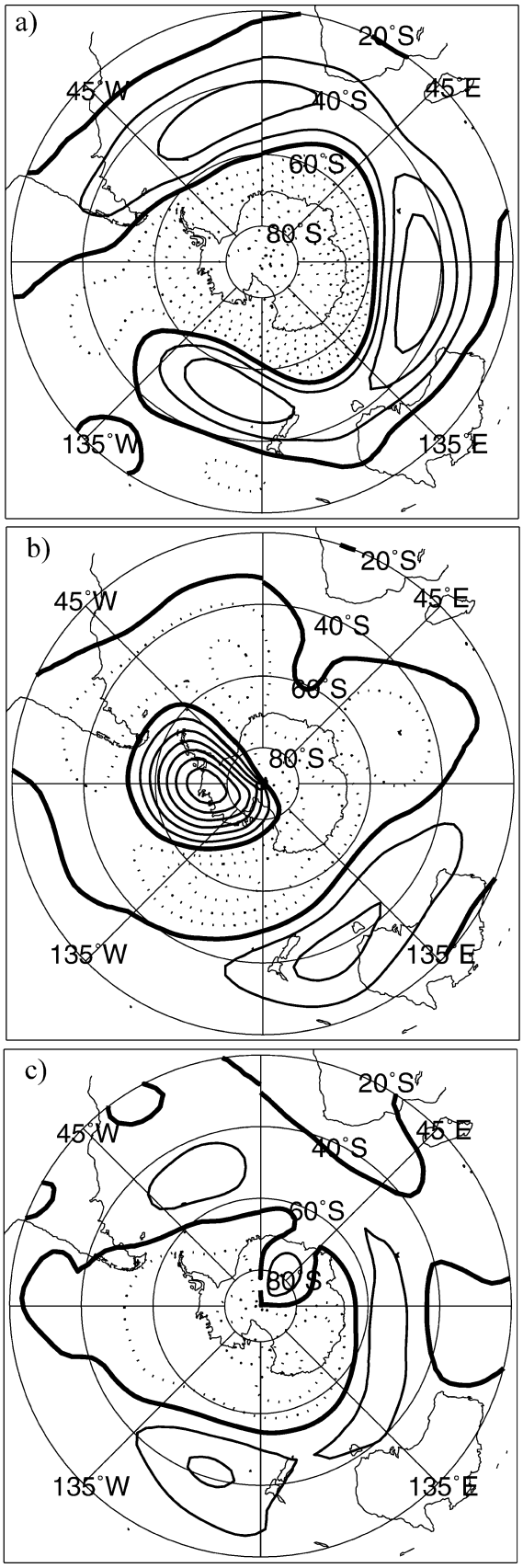

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FIG. 9. Regression of *Z*500 data upon the (a) first, (b) second, and (c) third normalized T_{IR} expansion coefficients shown in Fig. 4. Units, contours as in Fig. 6.

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