NOTES AND CORRESPONDENCE

Do Models and Observations Disagree on the Rainfall Response to Global Warming?

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ABSTRACT

Recently analyzed satellite-derived global precipitation datasets from 1987 to 2006 indicate an increase in global-mean precipitation of 1.1%–1.4% decade⁻¹. This trend corresponds to a hydrological sensitivity (HS) of 7% K⁻¹ of global warming, which is close to the Clausius–Clapeyron (CC) rate expected from the increase in saturation water vapor pressure with temperature. Analysis of two available global ocean evaporation datasets confirms this observed intensification of the atmospheric water cycle. The observed hydrological sensitivity over the past 20-yr period is higher by a factor of 5 than the average HS of 1.4% K⁻¹ simulated in state-of-the-art coupled atmosphere–ocean climate models for the twentieth and twenty-first centuries. However, the analysis shows that the interdecadal variability in HS in the models is high—in particular in the twentieth-century runs, which are forced by both increasing greenhouse gas (GHG) and tropospheric aerosol concentrations. About 12% of the 20-yr time intervals of eight twentieth-century climate simulations from the third phase of the Coupled Model Intercomparison Project (CMIP3) have an HS magnitude greater than the CC rate of 6.5% K⁻¹. The analysis further indicates different HS characteristics for GHG and tropospheric aerosol forcing agents. Aerosol-forced HS is a factor of 2 greater, on average, and the interdecadal variability is significantly larger, with about 23% of the 20-yr sensitivities being above the CC rate. By thermodynamically constraining global precipitation changes, it is shown that such changes are linearly related to the difference in the radiative imbalance at the top of the atmosphere (TOA) and the surface (i.e., the atmospheric radiative energy imbalance). The strength of this relationship is controlled by the modified Bowen ratio (here, global sensible heat flux change divided by latent heat flux change). Hydrological sensitivity to aerosols is greater than the sensitivity to GHG because the former have a stronger effect on the shortwave transmissivity of the atmosphere, and thus produce a larger change in the atmospheric radiative energy imbalance. It is found that the observed global precipitation increase of 13 mm yr⁻¹ decade⁻¹ from 1987 to 2006 would require a trend of the atmospheric radiative imbalance (difference between the TOA and the surface) of 0.7 W m⁻² decade⁻¹. The recovery from the El Chichón and Mount Pinatubo volcanic aerosol injections in 1982 and 1991, the satellite-observed reductions in cloudiness during the phase of increasing ENSO events in the 1990s, and presumably the observed reduction of anthropogenic aerosol concentrations could have caused such a radiative imbalance trend over the past 20 years. Observational evidence, however, is currently inconclusive, and it will require more detailed investigations and longer satellite time series to answer this question.

1. Introduction

Changes in the global hydrological cycle associated with greenhouse gas–induced warming are one of the most important aspects of anthropogenic climate change. The common assumption is that global precipitation will increase in a warmer world as a result of the strong temperature dependence of the water vapor saturation pressure $e_s$, as given by the Clausius–Clapeyron (CC) equation. For a global-mean surface temperature of 288 K (15°C), the CC equation predicts that $e_s$ will increase by 6.5% K⁻¹ of surface warming.

Climate models (e.g., Allen and Ingram 2002) confirm an increase in the total amount of water vapor in the atmosphere at approximately the CC rate. Observations (e.g., Trenberth et al. 2005) indicate an increase of 1.3% ± 0.3% decade⁻¹ from 1988 to 2001. With the earth’s surface warming of 0.2 K decade⁻¹ over the
same period (Trenberth et al. 2007), a rate of total water vapor increase of indeed 7% K^{-1} is observed. Climate models, however, suggest that global-mean \( P \) increases at a much slower rate of about 1\%-3\% K^{-1} of global warming (e.g., Held and Soden 2006). This discrepancy has been explained by either the thermodynamic energy constraints on the water cycle, which require global-mean \( P \) to change in accordance with the difference between the top of the atmosphere (TOA) and the surface radiative balance (see Mitchell et al. 1987; Boer 1993; Allen and Ingram 2002; Feichter et al. 2004), or by the dynamics of the Hadley circulation and extratropical moisture transport (Held and Soden 2006; Lorenz and DeWeaver 2007), which also dampen the global-mean \( P \) response to warming.

However, in a recent paper Wentz et al. (2007, hereinafter W07) state that over the last two decades global precipitation increased at close to the CC rate. This finding is based on satellite-derived global-mean \( P \) from Special Sensor Microwave Imager (SSM/I; version 6) that increased by 13.2 \pm 4.8 \text{ mm yr}^{-1} \text{ decade}^{-1} (1.4\% \pm 0.5\% \text{ decade}^{-1}) over the period from July 1987 through August 2006. The latest analysis of global-mean \( P \) from the Global Precipitation Climatology Project (GPCP; Adler et al. 2003; Gu et al. 2007) results in a smaller long-term trend of 3.8 \text{ mm yr}^{-1} \text{ decade}^{-1} for the extended time period from 1979 to 2006 (Adler et al. 2008). In general, Yin et al. (2004) urge caution when using GPCP and Climate Prediction Center Merged Analysis of Precipitation (CMAP; Xie and Arkin 1998) precipitation products for trend analysis prior to SSM/I in 1987—in particular over the oceans because of input data changes and atoll sampling in CMAP. The Intergovernmental Panel on Climate Change Fourth Assessment Report (IPCC AR4; Solomon et al. 2007) concludes that from 1979 to 2005 global precipitation trends range from −16 to +13 \text{ mm yr}^{-1} \text{ decade}^{-1} for various datasets but that none of the trends are significant.

In accordance with other model studies, W07 found that decadal trends in \( P \) in the second Atmospheric Model Intercomparison Project (AMIP-II) model climate simulations (available online at http://www-pcmdi.llnl.gov/projects/amip) are smaller by a factor of 2–3 relative to the W07 observations for the 1987–2001 time period, even though these simulations are forced with observed sea surface temperatures. Allan and Soden (2007) further report of large discrepancies between observed (GPCP, CMAP, and SSM/I) and simulated precipitation trends of fully coupled [phase 3 of the Coupled Model Intercomparison Project (CMIP3)] model climate models in the ascending and descending branches of the tropical Hadley circulation. Increasing rainfall in the rising branch and decreasing trends in the descending regimes are detected in the observations for 1979–2006 [see also Chen et al. (2002), who confirm this feature], but these trends are not fully reproduced by the models. Both studies imply that the model–observation discrepancies in precipitation must be due to model parameterization errors and/or errors in the satellite data.

In this note, we focus on evaluating model performances on interdecadal time scales and argue that variability in climate-forcing agents such as aerosols and greenhouse gases (GHGs) on these time scales can play an important role in determining variability in observed and simulated global precipitation. We show that the strong increase in precipitation observed in the 1990s is not outside the CMIP3 models’ range of interdecadal variability. Last, we discuss how the different effects of aerosols and GHGs on the TOA and surface radiation may explain variations in global precipitation, and we describe what role the Bowen ratio plays (the ratio of surface sensible heat flux to latent heat flux).

2. Data

Here we use two monthly mean precipitation products for calculating linear trends for the period from July 1987 to August 2006 and for the model comparisons: W07’s merged product of ocean SSM/I and land GPCP data, and the GPCP product (Adler et al. 2003) for land and ocean precipitation. These global datasets have been analyzed before and are described in the introduction. Estimates of global ocean evaporation from synthesized surface meteorological conditions obtained from satellite remote sensing and atmospheric model reanalysis are also used. The evaporation products are the Objectively Analyzed Air–Sea Fluxes (OAFlux) for the global ice-free oceans by Yu and Weller (2007) and the Hamburg Ocean–Atmosphere Parameters and Fluxes from Satellite Data, version 3, (HOAPS3) by Andersson et al. (2007) for the oceans equatorward of 80° latitude. For both datasets, evaporation \( E \) at the ocean surface is calculated from the aerodynamic bulk formula based on Fairall et al. (2003) with wind speed, atmospheric specific humidity, and saturation specific humidity at sea surface as input data. The main difference of the datasets stems from the choice of input datasets and the preparation of the input data. OAFlux uses National Centers for Environmental Prediction and European Centre for Medium-Range Weather Forecasts reanalysis products and SSM/I, Quick Scatterometer, Advanced Very High Resolution Radiometer, Tropical Rainfall Measuring Mission Microwave Imager, and Advanced Microwave Scanning Radiometer for Earth Observing System satellite retrievals, whereas HOAPS3 relies only on SSM/I satellite
retrievals for the input parameters. For this study, we analyze \( E \) data for the time period from January 1987 to December 2004.

The model data used in our analysis are the global gridded monthly fields of \( P \) and surface air temperature (SAT) obtained for eight CMIP3 models (Geophysical Fluid Dynamics Laboratory Climate Model, version 2.0 (GFDL CM2.0); GFDL Climate Model, version 2.1 (GFDL CM2.1); Goddard Institute for Space Studies Model E-H (GISS-EH); Institute of Numerical Mathematics Coupled Model, version 3.0 (INM-CM3.0); Model for Interdisciplinary Research on Climate 3.2, high-resolution version [MIROC3.2(hires)]; MIROC3.2, medium-resolution version [MIROC(medres)]; Meteorological Institute of the University of Bonn “ECHO-G” Model (MIUBECHOG), and National Center for Atmospheric Research Community Climate System Model, version 3 (CCSM3)) for two different experiments, the “climate of the twentieth century” (20C3M) experiment and the Special Report on Emissions Scenarios A1B experiment. Both experiments were forced with observed (20C3M) or projected (A1B) changes in anthropogenic well-mixed greenhouse gases and aerosols. The 20C3M experiment also contains observationally based estimates of natural forcing resulting from changes in solar irradiance and volcanic activity. Because the initial conditions for the A1B experiment were taken from the end of the 20C3M experiment, continuous time series of \( P \) and SAT can be created for the entire twentieth and twenty-first centuries. Unlike the 20C3M and A1B experiments analyzed in this study, the AMIP-II experiments used by W07 did not apply any time-varying changes in atmospheric GHG, stratospheric and tropospheric aerosol concentrations, or solar irradiances. Yu and Weller (2007) conclude that these temporally variable forcings may affect moisture transport and evaporation. Thus, analysis of the 20C3M and A1B simulations provides an indication of whether modeled \( P \) changes during the past 20 years are sensitive to the specification of these time-varying forcings.

For the second part of the study, we analyze special simulations performed with the atmosphere model of the National Aeronautics and Space Administration (NASA) Goddard Institute for Space Studies coupled to the dynamic ocean model of Russell (GISS_ER). This model configuration was used for the IPCC AR4 (Solomon et al. 2007) and is described in detail by Hansen et al. (2007, hereinafter H07). The GISS_ER model was run with the IPCC AR4 transient forcings from 1880 to 2003. Different experiments were carried out with either all forcings applied simultaneously in the model or with each forcing applied individually (H07). The three experiments analyzed here are the “tropospheric aerosols only” experiment (trAerosols), the “well-mixed greenhouse gases only” experiment (GHG), and the experiment with all forcings applied at once. We show model results based on the multimember ensemble means from each of these experiments. Time-dependent tropospheric aerosols in the GISS_ER model are sulfate, black carbon, organic carbon, and nitrate. From 1880 onward, tropospheric aerosol optical depths increase nonlinearly—in particular, after 1950—similar to well-mixed greenhouse gases. After 1990 the tropospheric aerosol forcing is kept constant, whereas greenhouse gases continue to increase. Details of temporal and spatial distributions of these aerosols can be found in Koch (2001) and Hansen et al. (2005b). The indirect effect of aerosols on cloud coverage is parameterized as dependence of cloud coverage on the logarithm of concentrations of soluble aerosols. The total aerosol (effective) forcing, including the indirect effect on cloud coverage, is \( F_c = -1.37 \text{ W m}^{-2} \) from 1880 to 2003. Well-mixed greenhouse gases in the model are carbon dioxide, methane, nitrous oxide, chlorofluorocarbons, and other trace gases. H07 report a GHG forcing of \( F_\text{c} = 2.72 \text{ W m}^{-2} \) for 1880–2003, which is 2 times the magnitude of the aerosol forcing. The individual model runs for these and a variety of other forcing agents were available online at the time of writing (http://data.giss.nasa.gov/modelE/transient/climsim.html).

3. Interdecadal variability of global precipitation, evaporation, and temperature

Figure 1 shows the linear trends in global mean \( P \) and SAT for the eight models for the period from July 1987 through August 2006. Also shown are the observed \( P \) and SAT trends (dark gray bars) during this time period from SSM/I datasets (see W07) and the GPCP global precipitation trend (light gray shading) over the same period, which is slightly lower. We do not show the global \( P \) trend of \(-0.138\%\) decade\(^{-1}\) of the CMAP dataset for the reasons discussed in the introduction. The figure also includes the GISS_ER trAerosol and GHG experiments, as well as the “post Mount Pinatubo” experiment, which represents the GISS_ER simulations, including all forcings, sampled here for the 12-yr period immediately following the Mount Pinatubo volcanic eruption in June of 1991.

It is obvious in Fig. 1 that all CMIP3 models underestimate the observed \( P \) increase of \(1.40\% \pm 0.50\%\) decade\(^{-1}\) for SSM/I and \(1.14\% \pm 0.65\%\) decade\(^{-1}\) for GPCP during the past 20 years despite the fact that six of the eight models simulate more surface warming than was observed. (Note the large uncertainty ranges of the observed \( P \) trends.) We assessed the statistical significance
of the difference between observed and modeled \( P \) trends following Santer et al. (2000a,b) and found the difference to be significant at the 5% level for all models and the SSM/I dataset, but for only two of the eight models and the GPCP dataset. The ratio of the global-mean \( P \) trend to the global-mean SAT trend, referred to herein as the hydrological sensitivity (HS; e.g., Liepert et al. 2004), varies between 0.71% and 4.0% K\(^{-1}\) for the eight models, with a multimodel mean of 1.9% K\(^{-1}\). The average HS is clearly smaller than the CC rate, and this result is in line with earlier work. Despite the large uncertainties of precipitation observations, several authors recognized the high interdecadal variability of observed global precipitation resulting from internal climate variability such as ENSO or external forcings such as major volcanic eruptions (Yin et al. 2004; Gu et al. 2007; Trenberth and Dai 2007; Trenberth et al. 2007). Gu et al. for example, estimate a contribution as high as a 5% reduction of tropical rainfall following several years after volcanic eruptions (Mount Pinatubo and El Chichón).

Figure 2 shows the global-mean anomalies of ocean surface evaporation from the OAFlux and HOAPS3 datasets for the slightly shorter time period from January 1987 to December 2004. The seasonality of \( E \) as well as \( P \) datasets was removed before the gridbox monthly mean anomalies were calculated. Evaporation anomalies over sea ice grid boxes in OAFlux and...
poleward of 80° in HOAPS3 are generally small relative to lower latitudes and were set to zero for this analysis. The 12-month running means (thick lines) indicate an upward tendency of the global ocean evaporation after a drop in 1991, which coincides with the Mount Pinatubo volcanic eruption in June of 1991. Evaporation drops dramatically in the HOAPS3 dataset following the Mount Pinatubo eruption but drops only slightly in OAFlux. Even after the Mount Pinatubo recovery, ocean evaporation after 2000 clearly exceeds the evaporation of the years before Mount Pinatubo in both datasets. The El Chichón volcanic eruption in March 1982 may have affected the beginning of the data as suggested by Gu et al. (2007). The main differences in $E$ between the datasets may be due to the inclusion of reanalysis data in OAFlux, which do not account for the volcanic aerosol forcing and its full effects on dynamics and surface wind speed. HOAPS3 evaporation estimates only rely on SSM/I satellite retrievals of wind, humidity, and temperature fields, whereas OAFlux synthesizes reanalysis and satellite products of various sources. A more detailed investigation is needed to clarify these differences. We estimate a long-term increase in ocean evaporation of $18 \pm 9$ mm yr$^{-1}$ decade$^{-1}$ and $51 \pm 39$ mm yr$^{-1}$ decade$^{-1}$ from 1987 to 2004 for OAFlux and HOAPS3 based on a linear trend analysis. Note that the linear trend mainly reflects the low-frequency change and not the Mount Pinatubo dip particularly seen in HOAPS3. If no changes in evapotranspiration from land areas are assumed, the global-mean trends of evaporation solely based on ocean changes would be $16 \pm 8$ mm yr$^{-1}$ decade$^{-1}$ and $45 \pm 35$ mm yr$^{-1}$ decade$^{-1}$ for OAFlux and HOAPS3, respectively. (We note that ocean evaporation is about a factor of 6 larger than evapotranspiration from land.) Thus, OAFlux evaporation changes without land contribution compare favorably to the global-mean $P$ increase of $13 \pm 5$ mm yr$^{-1}$ decade$^{-1}$ from 1987 to 2006 (W07), whereas HOAPS3 changes would likely require reductions of evapotranspiration from land to match the observed global $P$ trends. However, both $E$ datasets are consistent with the observed $P$ increase to within their respective uncertainty ranges.

Apart from the uncertainties in the global observational data, the discrepancies between global-mean modeled and observed $P$ changes as shown in Fig. 1 during the past 20 years lead one to question how well the interdecadal
variability of global-mean $P$ or $E$ is reproduced in models and how this interdecadal variability will change in the future in response to anthropogenic forcing.

Here we examine the interdecadal variations in $P$ in the CMIP3 climate models to evaluate the likelihood of a 7% global $P$ increase per kelvin, as recently observed in the 20-yr SSM/I merged satellite product. Linear trends in global-mean $P$ for overlapping 20-yr periods in the twentieth and twenty-first centuries are calculated. Figure 3 shows the combined (i.e., for all eight models) distributions of $P$ trends along with distributions of global-mean SAT trends and the hydrological sensitivities based on these trends. The mean 20-yr global $P$ and SAT trends are 0.076% decade$^{-1}$ and 0.093 K decade$^{-1}$ in the twentieth century and increase considerably to 0.41% decade$^{-1}$ and 0.28 K decade$^{-1}$ in the twenty-first century. For comparison, the multimodel mean global $P$ and SAT trends for the period from July 1987 through August 2006 (see Fig. 1) are 0.55% decade$^{-1}$ and 0.29 K decade$^{-1}$, respectively. Thus, this recent 20-yr period in the models is anomalous in terms of its relatively rapid rate of $P$ increase when compared with the mean rate of $P$ increase during the twenty-first century, even though the corresponding rates of surface warming are very similar. It is evident from Fig. 3 that several 20-yr periods in both the twentieth and twenty-first centuries are characterized by $P$ trends falling within the uncertainty range of the observed $P$ trends. Model-simulated trends are as large as the best estimate of 1.14% decade$^{-1}$ of GPCP but are not as large as the best estimate of 1.40% decade$^{-1}$ of the SSM/I precipitation trend.
The above results indicate that 20-yr increases in global $P$ close to the observed increases are not outside the models’ range of interdecadal variability, albeit the simulated mean trend is considerably smaller. Such relatively large $P$ increases occur in the models because of some combination of internal climate variability, interdecadal variability in forcing, and/or climate response to forcing. The latter will be investigated below. Figure 3 also indicates that several 20-yr periods of decreasing global $P$ are predicted to occur by the models in the twenty-first century despite the fact that all 20-yr periods in the century show increasing SAT. That global $P$ can decrease in a warming world has been suggested previously (Liepert et al. 2004) as a consequence of anthropogenic aerosol emission increases.

The distributions of HS shown in Fig. 3 reveal notable differences between the twentieth and twenty-first centuries. Although the median 20-yr HS is the same in both centuries, being about 1.4% K$^{-1}$, the spread of the HS distribution is significantly greater during the twentieth century. (Note that the median rather than the arithmetic mean is used as an indication of the average 20-yr HS because of the presence of large-magnitude outliers in the distribution.) Approximately 12% of the 20-yr intervals in the twentieth century have an HS that is greater than the increase in moisture holding capacity (CC rate of 6.5% K$^{-1}$). These large-magnitude HS values are generally associated with very small magnitude trends in global-mean SAT. For example, for the periods with an HS larger in magnitude than the CC rate, the mean absolute value of the 20-yr SAT trend is only 0.017 K decade$^{-1}$. In the twenty-first century, as global warming accelerates, these small-magnitude SAT trends are reduced considerably in number and the HS distribution narrows (see Fig. 3). The entire twenty-first-century HS distribution is within plus or minus the CC rate.

4. Hydrological sensitivity to GHG and tropospheric aerosol forcings

We expect that differences in the width of the HS distributions between the twentieth and the twenty-first century arise partly as a result of differences in the forcings and their climate feedbacks. The A1B (twenty-first century) scenario is dominated by increasing well-mixed greenhouse gases, whereas in the twentieth-century...
forcing scenario increases in GHGs are accompanied by other forcings, predominantly rising aerosol concentrations (H07). Different HS to aerosol and greenhouse gas forcings exemplifies the different $P$ responses to shortwave (SW) and longwave (LW) forcings. Gillett et al. (2004) detected the influence of volcanic aerosols on observed global precipitation and concluded that “shortwave forcings exert a larger influence on precipitation than longwave forcings.”

Here we assess the difference in HS to aerosol and greenhouse gas forcing agents. The trend analysis for 20-yr overlapping periods is repeated with results from the GISS_ER GHG and trAerosol experiments. As expected, the average global $P$ trend of 0.31% decade$^{-1}$ and SAT trend of 0.16 K decade$^{-1}$ for the GHG experiment are smaller than the corresponding trends for the twenty-first-century scenario (A1B) but are substantially higher than the twentieth-century 20C3M simulations that include all forcings (see section 2). In contrast to GHG, increasing tropospheric aerosols in the trAerosol scenario exert a small mean cooling trend of $-0.085$ K decade$^{-1}$ and a mean decrease in global $P$ of $-0.39$% decade$^{-1}$. The median hydrological sensitivity of 4.1% K$^{-1}$ (mean of 5.1% K$^{-1}$) of the trAerosol experiment is about a factor of 2 larger than the median HS of the GHG experiment, which is 1.9% K$^{-1}$ (mean of 1.7% K$^{-1}$). The observed HS value of 7% K$^{-1}$ from the SSM/I data for the period of 1987-2006 is well within the standard deviation of $\pm 2.6\%$ K$^{-1}$ of the mean HS of trAerosol. The 20-yr HS distributions for the GHG and trAerosol simulations are shown in Fig. 4. As expected, the GHG distribution is very similar to the twenty-first-century HS distribution in Fig. 3, which is dominated by greenhouse gas warming. Most of the hydrological sensitivities fall between 0% and 3% K$^{-1}$, which is much smaller than the CC rate. Note that, similar to the twenty-first-century distributions (Fig. 3), about 6% of the 20-yr global $P$ trends in GHG are negative despite positive SAT trends for these periods, which indicates large interdecadal variability. This is not the case for trAerosol, for which the entire HS distribution is positive and substantially broader than GHG, with approximately 23% of the 20-yr sensitivities above the CC rate. Thus, the analysis of hydrological sensitivity and its interdecadal variability indicates different characteristics for different forcing agents.

5. Thermodynamic control of precipitation changes

To close the atmospheric water budget $P$ must equal $E$ in a global, annual sense, and changes in precipitation $\delta P$ must equal changes in evaporation $\delta E$ at the surface. Further, precipitation changes $\delta P$ are accompanied by anomalous release of latent heat $\delta Q_{\text{lat}}$ resulting from phase transition. Hence, long-term global precipitation changes, which are discussed here, are linked to changes in the surface energy budget as follows:

$$L \delta P = L \delta E = \delta Q_{\text{lat}} = \delta R_{\text{surf}} - \delta Q_{\text{sen}} - \delta M. \quad (1)$$

Here, $L$ is the energy of phase transition, $\delta R_{\text{surf}}$ is the net radiative energy change at the surface, $\delta Q_{\text{sen}}$ is the sensible heat change, and $\delta M$ is the ocean and land heat uptake (plus a small contribution from changes in melting of snow and ice). We introduce a modified Bowen ratio $B^*$ that links the global-mean changes of sensible heating $\delta Q_{\text{sen}}$ and latent heating $\delta Q_{\text{lat}}$ as follows:

$$B^* = \frac{\delta Q_{\text{sen}}}{\delta Q_{\text{lat}}}. \quad (2)$$

Equation (1) can then be rewritten as

$$\delta P = \frac{\delta R_{\text{surf}} - \delta M}{L (1 + B^*)}. \quad (3)$$

On decadal time scales the ocean and land heat uptake $\delta M$ is equal to the radiative imbalance at the top of
the atmosphere $\delta R_{\text{TOA}}$ (Hansen et al. 2005a). Replacing $\delta M$ with $\delta R_{\text{TOA}}$ in Eq. (3) and dividing both sides of the equation by the global-mean surface air temperature change $\delta T$ leads to

$$L(1 + B^*) \frac{\delta P}{\delta T} = \frac{\delta R_{\text{ref}}}{\delta T} - \frac{\delta R_{\text{TOA}}}{\delta T}. \quad (4)$$

Thus, the modified energy balance in Eq. (4) links the hydrological sensitivity $\delta P/\delta T$ to the climate sensitivity, here the global surface air temperature response to the adjusted TOA radiative imbalance $\delta R_{\text{TOA}}/\delta T$. This means that HS is linearly related to the difference of the surface and TOA radiative imbalances, with a slope determined by the modified Bowen ratio $B^*$.

Figure 5 illustrates this linear behavior in the GISS_ER model simulations. Shown is the scatterplot of the global overlapping 20-yr trends of monthly precipitation anomalies for the GHG (red asterisks) and the trAerosol (blue asterisks) simulations against the trends of the differences of surface-minus-TOA net radiative flux monthly anomalies. The slope of the linear regression line through all of the data (i.e., GHG and trAerosol) is 12.7 mm yr$^{-1}$ (W m$^{-2}$)$^{-1}$. This slope corresponds to a $B^*$ of $-0.0065$, which means that 20-yr sensible heat flux trends are significantly smaller and anticorrelated to latent heat flux trends for the same time periods. Recalculating the slope for each experiment separately, we find modified Bowen ratios of $-0.30$ for GHG and $+0.40$ for trAerosol. Hence, for the GHG experiment the sensible heat flux trends tend to be one-third in magnitude and anticorrelated to the latent heat flux trends, whereas in the trAerosol case sensible and latent heat flux trends tend to be correlated.

6. Summary and conclusions

A new satellite-derived analysis of global precipitation (W07) detected an increase of global-mean $P$ of $13.2 \pm 4.8$ mm yr$^{-1}$ decade$^{-1}$ ($1.4 \pm 0.5$% decade$^{-1}$) over the period from July 1987 through August 2006. With a surface air temperature increase of 0.2 K decade$^{-1}$, this trend amounts to a hydrological sensitivity of 7% K$^{-1}$, which is close to the Clausius–Clapeyron rate that describes the dependence of water vapor saturation pressure on temperature. Our analysis of two newly available ocean evaporation datasets for January 1987–December 2004 indicates an increase in global-mean $E$ that is consistent with the W07 global-mean $P$ trend within the limits of data uncertainty. The linear trends in global $E$ are significant at the 95% confidence level with $16 \pm 8$ mm yr$^{-1}$ decade$^{-1}$ for the OAFlux estimates and $45 \pm 35$ mm yr$^{-1}$ decade$^{-1}$ for the HOAPS3 estimates. A strong drop in $E$ following the Mount Pinatubo volcanic eruption in June 1991 is a prominent feature in HOAPS3 but not in the OAFlux dataset. Differences in the choice of input data may cause this discrepancy, and further, more detailed investigation is needed.

Our analysis confirms the W07 result that state-of-the-art climate models underestimate the observed 7% K$^{-1}$ hydrological sensitivity during 1987–2006. We show in this study, nonetheless, that CMIP3 coupled atmosphere–ocean climate models can produce an HS this large on 20-yr time scales. About 12% of 20-yr time periods of the twentieth-century climate simulations are characterized by HS equal to or larger than the CC rate. However, the average HS for all 20-yr periods in this century (and in the twenty-first-century simulations of the CMIP3 models) is much smaller, being about 1.4% K$^{-1}$. (Note that 6% of the 20-yr intervals in the twenty-first-century simulations exhibit global warming concurrent with declining precipitation.) Thus, the global $P$ change during a given 20-yr period may not be representative of the changes that will occur on longer time scales.

The main difference between the twentieth and twenty-first century in the models is the dominance of the greenhouse gas forcing over the tropospheric aerosol forcing in the current century. The analysis of special GISS_ER climate runs shows that HS of simulations with only twentieth-century tropospheric aerosol forcing (trAerosol) can be up to 5 times as large as the HS of simulations forced by twentieth-century GHGs only. The median 20-yr HS is 4% K$^{-1}$ for the trAerosol experiment and 2% K$^{-1}$ for the GHG experiment. The HS distribution of the trAerosol experiment is substantially broader, with 23% of the 20-yr time intervals characterized by an HS above the CC rate.

We further investigated the relationship between hydrological sensitivity and climate sensitivity (here the ratio of TOA radiative imbalance to global-mean surface air temperature change over a multiannual time period). In a global, long-term mean sense, hydrological sensitivity can be related to the energy imbalance of the atmosphere. We show that long-term variations in surface-minus-TOA radiative imbalance (atmospheric radiative energy imbalance) correlate linearly with precipitation variations. The strength of this linear relationship (i.e., the slope) is determined by the modified Bowen ratio $B^*$, which is here defined as sensible heat flux changes divided by latent heat flux changes at the surface. (Note that the Bowen ratio is generally defined as average sensible heat flux divided by average latent heat flux.)

We show that the GISS, ER model reproduces this analytically derived linear relationship very well. Latent heat changes and sensible heat changes in the model are anticorrelated in the GHG experiment and correlated in the trAerosol experiment ($B^*$ of $-0.3$ and $0.4$, respectively).
respectively). When data from both experiments are combined, latent heat flux changes are significantly larger than and anticorrelated with sensible heat flux changes.

Equation (4) states that the hydrological sensitivity is controlled by changes in the atmospheric radiative energy imbalance. This energy imbalance exists largely because the atmosphere is relatively transparent to SW radiation, and thus atmospheric SW absorption is smaller than the net LW emission from the atmosphere. Therefore, changes in the atmospheric radiative energy imbalance should be proportional to changes in the SW transmissivity of the atmosphere. Because clouds and aerosols significantly affect the SW transmissivity by reflecting and absorbing incoming SW radiation (e.g., Liepert 2002; Liepert and Tegen 2002; Kim and Ramanathan 2008), we expect them generally to produce a stronger HS than greenhouse gases, which have comparatively little effect on SW transmission. This is supported by our analysis of the GISS_ER GHG and trAerosol experiments. In the latter experiment, changes in both surface sensible heat flux and latent heat flux act to counter the changes in atmospheric radiative energy imbalance (i.e., positive $B^*$), whereas in GHG sensible and latent heat flux changes oppose one another (negative $B^*$).

Based on our model analysis, we find that an increase in the difference between global surface and TOA net radiation of 0.7 W m$^{-2}$ decade$^{-1}$ would be needed to account for the observed precipitation increase of 13 mm yr$^{-1}$ decade$^{-1}$ from 1987 to 2006. We investigated global-mean radiative flux anomalies for the time period from July 1987 to December 2005 using the Surface Radiation Budget (SRB) dataset from the NASA Langley Research Center Atmospheric Sciences Data Center (available online at http://eosweb.larc.nasa.gov/PRDOCS/srb/table_srb.html) and found that the difference between surface and TOA net radiation increased by 0.8 W m$^{-2}$ decade$^{-1}$ [see also Zhang et al. (2007) for an alternative dataset]. However, even for the deseasonalized monthly anomalies, the standard deviation of this estimate is about 2 W m$^{-2}$ decade$^{-1}$, signifying large uncertainty in the radiative flux measurements and thus precluding any definitive conclusion about whether the flux changes are consistent with the observed $P$ change.

The mid-1980s and mid-1990s saw a decrease in stratospheric aerosol injected by the eruptions of the volcanoes El Chichón (1982) and Mount Pinatubo (1991). Global reductions in anthropogenic aerosol optical depths (Mishchenko et al. 2007) and decreases in global cloudiness (Wielicki et al. 2002; Romanou et al. 2007) resulting from more frequent El Niño events were also recently observed by satellites. These cloud and aerosol changes induced a “brightening” trend that may have contributed to the rapid increase in global $P$ observed during the past two decades (see also Trenberth and Dai 2007). Detection of “dimming and brightening” trends, however, is still subject of ongoing investigation (see Hinkelmann et al. 2008, manuscript submitted to J. Geophys. Res.). We also suggest that more observational studies are needed to evaluate the relationship between hydrological sensitivity and the radiative forcing of climate, as we have done here using the GISS model.

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